

Stress Tensor

$$\underline{T}^n(\underline{x}) = \underline{\underline{\sigma}}(\underline{x}) \cdot \underline{n}$$

Stress tensor

The stress tensor depends on $\underline{x}$ alone.

To obtain the stress tensor at a point:

- choose three mutually perpendicular planes at that point,
- find tractions on those planes
- get the components of the tractions on the planes

Stress matrix as a representation of stress tensor

The state of stress at a point is completely defined by the **NINE** stress components acting on three mutually perpendicular planes (say $\underline{e}_1$ - $\underline{e}_2$ - $\underline{e}_3$ planes)

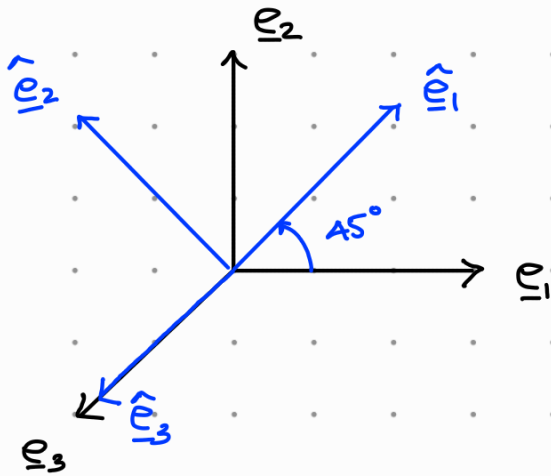
$$[\underline{\sigma}]_{(\underline{e}_1, \underline{e}_2, \underline{e}_3)} = \begin{bmatrix} \sigma_{11} & \tau_{21} & \tau_{31} \\ \tau_{12} & \sigma_{22} & \tau_{32} \\ \tau_{13} & \tau_{23} & \sigma_{33} \end{bmatrix}$$

$\underline{T}^1 \quad \underline{T}^2 \quad \underline{T}^3$

Suppose the stress matrix in $(e_1-e_2-e_3)$ coordinate system is given by

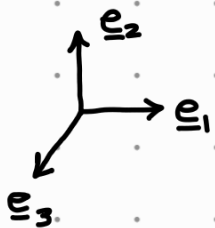
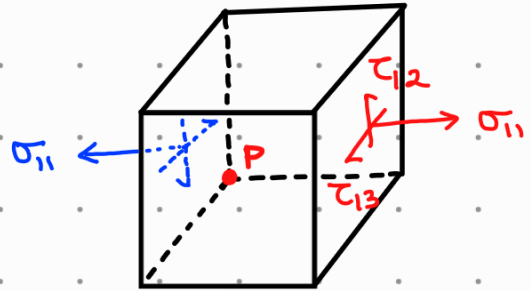
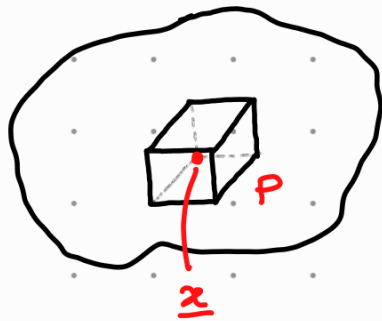
$$[\underline{\sigma}]_{(e_1-e_2-e_3)} = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Think of a new coordinate system obtained by rotation of $e_1-e_2-e_3$ coordinate system by 45° about e_3 . Find the stress matrix in the new coordinate system.



Cuboidal representation of stress tensor

To represent the stress tensor at a given point $\underline{x}$ in a body, we take an infinitesimally small cuboid

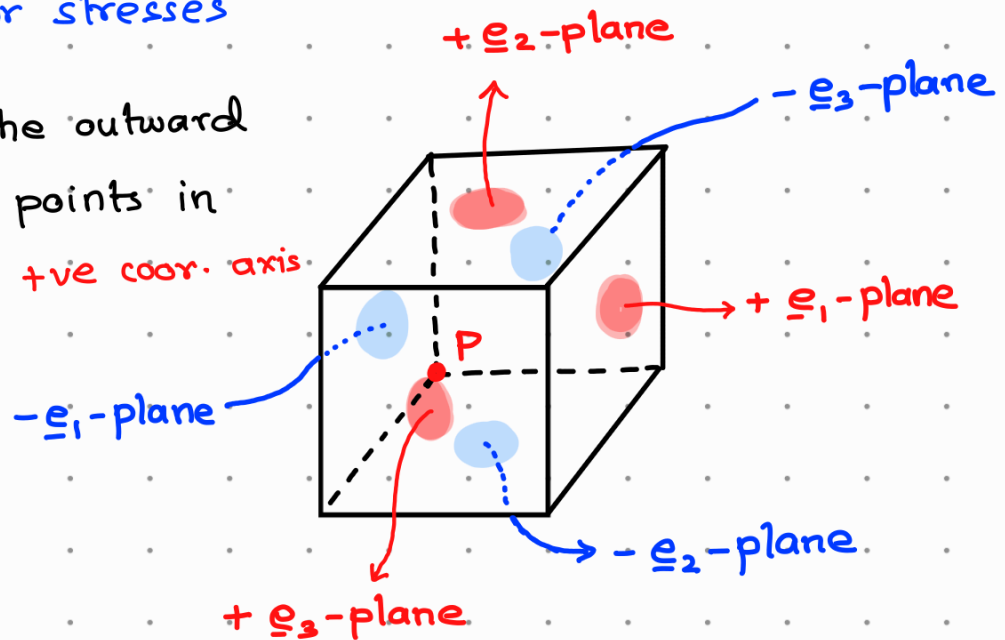


How do we define which stress is positive or negative?

→ We need a SIGN CONVENTION

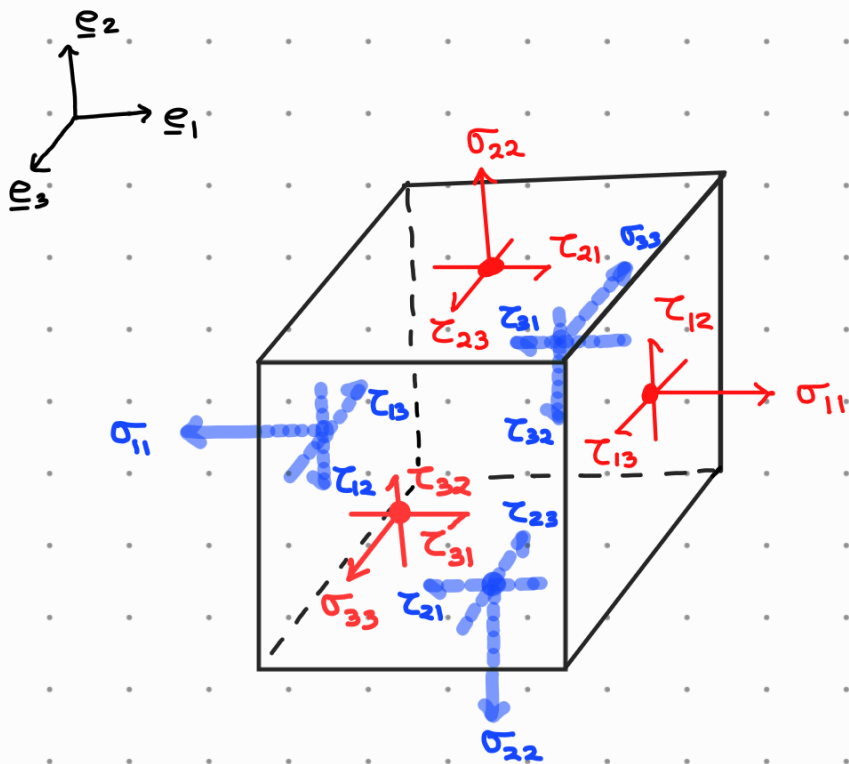
Sign Convention for stresses

- A face is **+ve** if the outward face normal $\underline{v}$ points in the direction of the **+ve** coord. axis.
- A face is **-ve** if the outward face normal points in the direction of the **-ve** coordinate axis.

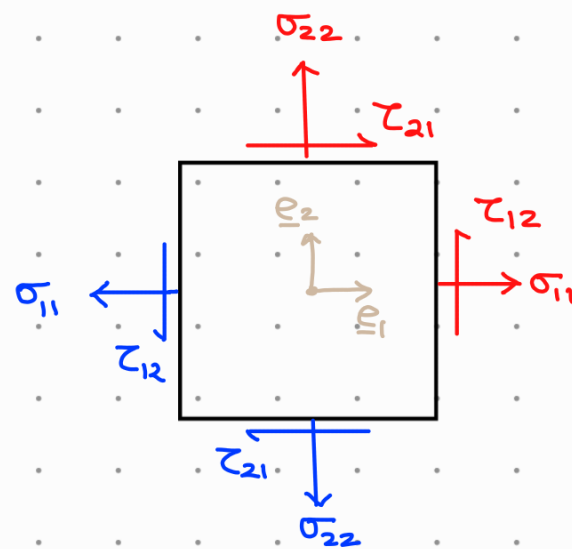


- The stress component is **+ve** when a positively directed force component acts on a positive face.
- The stress component is **+ve** when a negatively directed force component acts on a negative face.

- When a positively directed force component acts on a negative face or a negatively directed force component acts on a positive face, the stress component is **negative**



3D state of stress

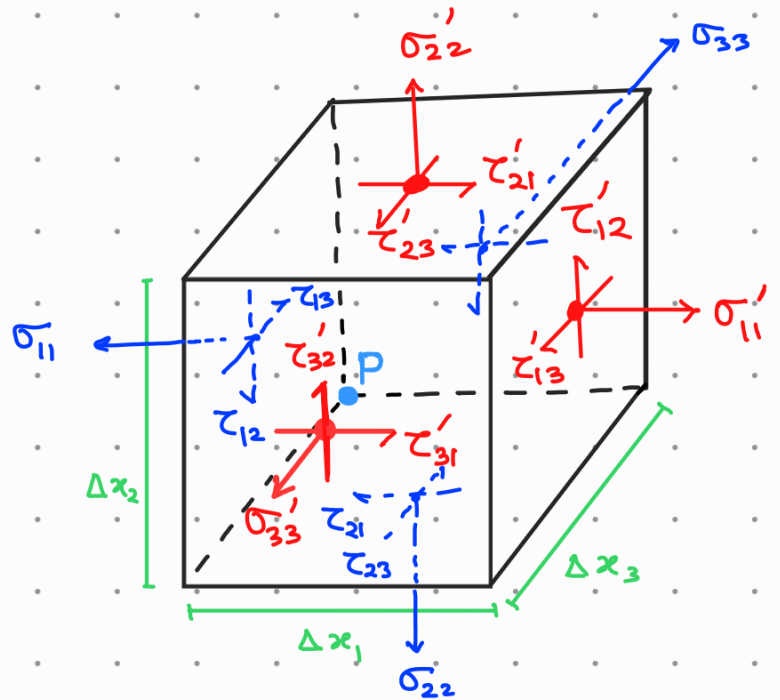
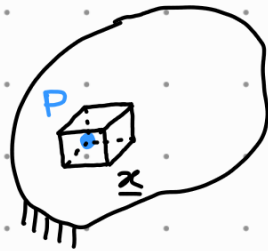
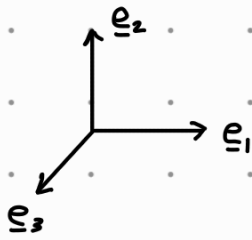


2D/plane state of stress

Certain remarks:

- The cuboid represents a very very small volume element from inside the body at the point P
- Note that all stress components shown are **positive**.
 - Positively directed force component act on +ve face
 - Negatively directed force component act on -ve face
- The stress components are assumed to be uniform over the face

Stress Equilibrium Equations in Cartesian Coordinates



The stress components on faces at distances Δx_1 , Δx_2 , and Δx_3 are different from the components at the faces passing through $P(x)$

How are σ_{11}' , τ_{12}' , τ_{13}' related to σ_{11} , τ_{12} , τ_{13} ?

Stress components on left face

$$\sigma_{11} = \sigma_{11}(x_1, x_2, x_3)$$

$$\tau_{12} = \tau_{12}(x_1, x_2, x_3)$$

$$\tau_{13} = \tau_{13}(x_1, x_2, x_3)$$

Stress components on right face

$$\sigma_{11}' = \sigma_{11}(x_1 + \Delta x_1, x_2, x_3)$$

$$\tau_{12}' = \tau_{12}(x_1 + \Delta x_1, x_2, x_3)$$

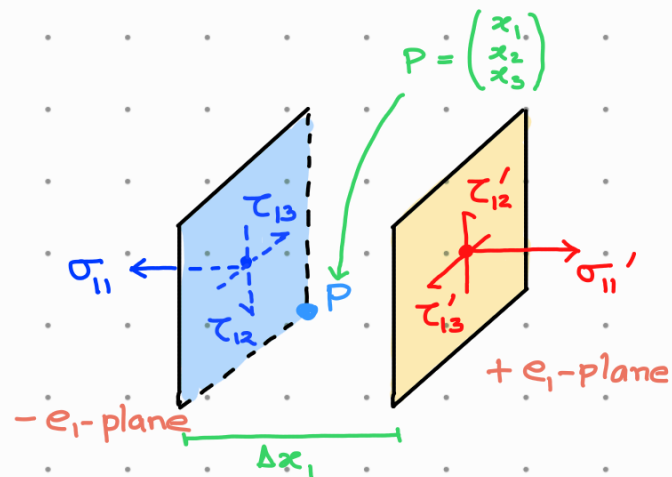
$$\tau_{13}' = \tau_{13}(x_1 + \Delta x_1, x_2, x_3)$$

Using Taylor series expansion,

$$\sigma_{11}' \approx \sigma_{11}(x_1, x_2, x_3) + \frac{\partial \sigma_{11}}{\partial x_1}(x_1, x_2, x_3) \Delta x_1$$

$$\tau_{12}' \approx \tau_{12}(x_1, x_2, x_3) + \frac{\partial \tau_{12}}{\partial x_1}(x_1, x_2, x_3) \Delta x_1$$

$$\tau_{13}' \approx \tau_{13}(x_1, x_2, x_3) + \frac{\partial \tau_{13}}{\partial x_1}(x_1, x_2, x_3) \Delta x_1$$



Face areas: $\Delta x_2 \Delta x_3$

Similarly, we can express relationships for other faces

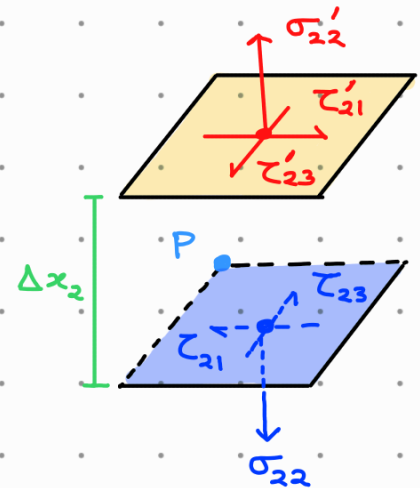
Face area: $\Delta x_1 \Delta x_3$

Stress components on bottom face

$$\sigma_{22} = \sigma_{22}(x_1, x_2, x_3)$$

$$\tau_{21} = \tau_{21}(x_1, x_2, x_3)$$

$$\tau_{23} = \tau_{23}(x_1, x_2, x_3)$$



Stress components on top face

$$\sigma_{22}' = \sigma_{22}(x_1, x_2 + \Delta x_2, x_3) \approx \sigma_{22}(\underline{x}) + \frac{\partial \sigma_{22}(\underline{x})}{\partial x_2} \Delta x_2$$

$$\tau_{21}' = \tau_{21}(x_1, x_2 + \Delta x_2, x_3) \approx \tau_{21}(\underline{x}) + \frac{\partial \tau_{21}(\underline{x})}{\partial x_2} \Delta x_2$$

$$\tau_{23}' = \tau_{23}(x_1, x_2 + \Delta x_2, x_3) \approx \tau_{23}(\underline{x}) + \frac{\partial \tau_{23}(\underline{x})}{\partial x_2} \Delta x_2$$

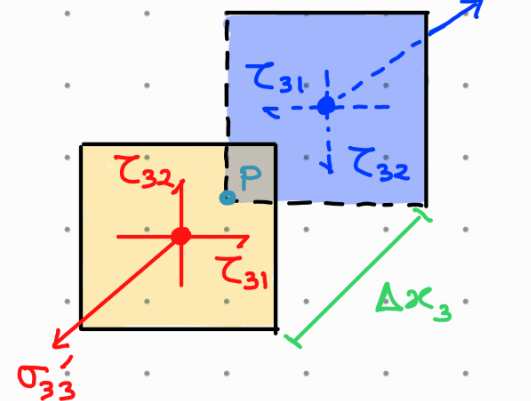
Stress components on back face

$$\sigma_{33} = \sigma_{33}(x_1, x_2, x_3)$$

$$\tau_{31} = \tau_{31}(x_1, x_2, x_3)$$

$$\tau_{32} = \tau_{32}(x_1, x_2, x_3)$$

Face area = $\Delta x_1 \Delta x_2$



Stress components on front face

$$\sigma_{33}' = \sigma_{33}(x_1, x_2, x_3 + \Delta x_3) \approx \sigma_{33}(\underline{x}) + \frac{\partial \sigma_{33}(\underline{x})}{\partial x_3} \Delta x_3$$

$$\tau_{31}' = \tau_{31}(x_1, x_2, x_3 + \Delta x_3) \approx \tau_{31}(\underline{x}) + \frac{\partial \tau_{31}(\underline{x})}{\partial x_3} \Delta x_3$$

$$\tau_{32}' = \tau_{32}(x_1, x_2, x_3 + \Delta x_3) \approx \tau_{32}(\underline{x}) + \frac{\partial \tau_{32}(\underline{x})}{\partial x_3} \Delta x_3$$

If a deformable body is in equilibrium, then **any isolated part of the body** must also be in equilibrium.

So, the cuboid must also be in equilibrium:

$$\sum \underline{M}_{\text{about any chosen pt}} = \underline{0} \Leftrightarrow \text{Balance of angular momentum}$$

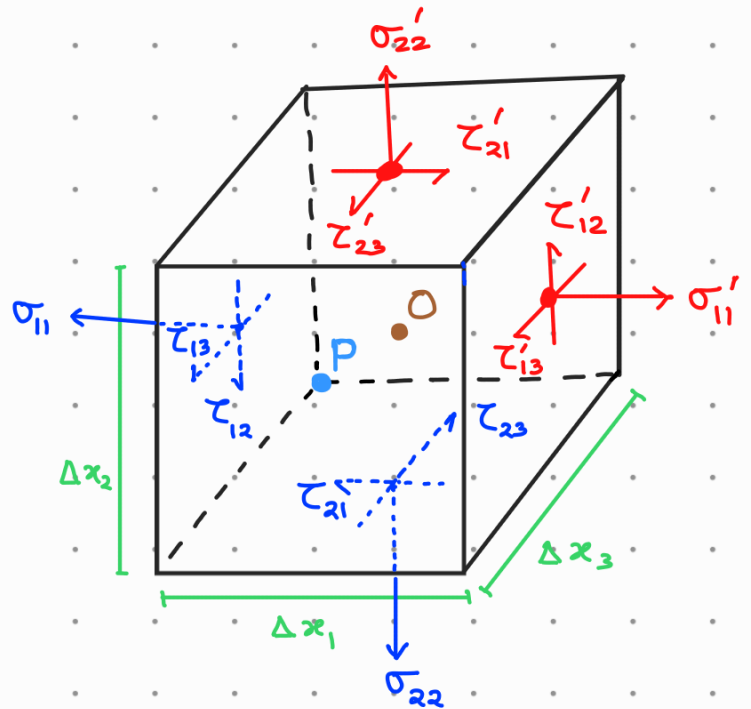
$$\checkmark \sum \underline{F} = \underline{0} \Leftrightarrow \text{Balance of linear momentum}$$

Let's consider the moment about the center of cuboid O

About the e_3 -axis

$$+ \curvearrowright \sum M_O = 0$$

$$\begin{aligned} \Rightarrow & (\tau_{12}' \Delta x_2 \Delta x_3) \frac{\Delta x_1}{2} - \\ & - (\tau_{21}' \Delta x_1 \Delta x_3) \frac{\Delta x_2}{2} - \\ & + (\tau_{12} \Delta x_2 \Delta x_3) \frac{\Delta x_1}{2} - \\ & - (\tau_{21} \Delta x_1 \Delta x_3) \frac{\Delta x_2}{2} = 0 \end{aligned}$$



$$\Rightarrow \left[\tau_{12} + \frac{\partial \tau_{12}}{\partial x_1} \Delta x_1 + \tau_{12} - \tau_{21} - \frac{\partial \tau_{21}}{\partial x_2} \Delta x_2 - \tau_{21} \right] \frac{\Delta x_1 \Delta x_2 \Delta x_3}{2} = 0$$

$$\Rightarrow 2\tau_{12} - 2\tau_{21} + \frac{\partial \tau_{12}}{\partial x_1} \Delta x_1 - \frac{\partial \tau_{21}}{\partial x_2} \Delta x_2 = 0$$

$$\text{As } \Delta x_1, \Delta x_2, \Delta x_3 \rightarrow 0$$

$$\boxed{\tau_{12} = \tau_{21}},$$

Likewise, you can get

$$\boxed{\tau_{13} = \tau_{31}},$$

$$\boxed{\tau_{23} = \tau_{32}}$$

Shear stress components on perpendicular faces are equal in mag.

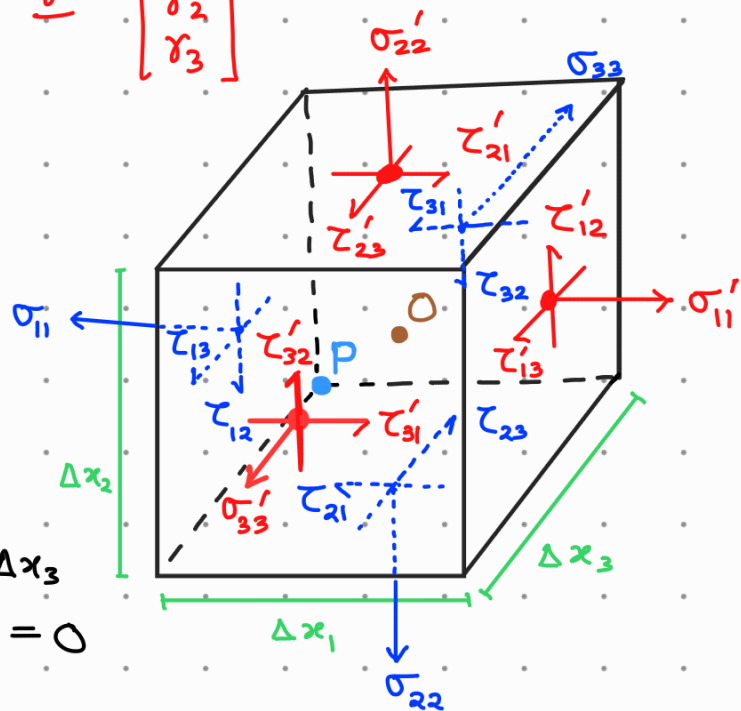
$$\rightarrow \sum F_x = 0$$

$$\begin{aligned} & \Rightarrow \sigma_{11}' \Delta x_2 \Delta x_3 - \sigma_{11} \Delta x_2 \Delta x_3 \\ & + \tau_{21}' \Delta x_1 \Delta x_3 - \tau_{21} \Delta x_1 \Delta x_3 \\ & + \tau_{31}' \Delta x_1 \Delta x_2 - \tau_{31} \Delta x_1 \Delta x_2 \\ & + \gamma_1 \Delta x_1 \Delta x_2 \Delta x_3 = 0 \end{aligned}$$

$$\Rightarrow \left[\frac{\partial \sigma_{11}}{\partial x_1} + \frac{\partial \tau_{21}}{\partial x_2} + \frac{\partial \tau_{31}}{\partial x_3} + \gamma_1 \right] \Delta x_1 \Delta x_2 \Delta x_3 = 0$$

$$\Rightarrow \frac{\partial \sigma_{11}}{\partial x_1} + \frac{\partial \tau_{21}}{\partial x_2} + \frac{\partial \tau_{31}}{\partial x_3} + \gamma_1 = 0$$

$$\underline{\gamma} = \begin{bmatrix} \gamma_1 \\ \gamma_2 \\ \gamma_3 \end{bmatrix}$$



Similarly,

$$+\uparrow \sum F_y = 0 \quad \Rightarrow \quad \frac{\partial \tau_{12}}{\partial x_1} + \frac{\partial \sigma_{22}}{\partial x_2} + \frac{\partial \tau_{32}}{\partial x_3} + \gamma_2 = 0$$

$$+\downarrow \sum F_z = 0 \quad \Rightarrow \quad \frac{\partial \tau_{13}}{\partial x_1} + \frac{\partial \tau_{23}}{\partial x_2} + \frac{\partial \sigma_{33}}{\partial x_3} + \gamma_3 = 0$$

Stress Equilibrium Relations

Three force eqⁿ + Three moment eqⁿ $\Rightarrow$ Total **SIX** relations.

Total **NINE** unknowns

$$\frac{\partial \sigma_{11}}{\partial x_1} + \frac{\partial \tau_{21}}{\partial x_2} + \frac{\partial \tau_{31}}{\partial x_3} + \gamma_1 = 0$$

$$\tau_{21} = \tau_{12}$$

$$\frac{\partial \tau_{12}}{\partial x_1} + \frac{\partial \sigma_{22}}{\partial x_2} + \frac{\partial \tau_{32}}{\partial x_3} + \gamma_2 = 0$$

$$\tau_{23} = \tau_{32}$$

$$\frac{\partial \tau_{13}}{\partial x_1} + \frac{\partial \tau_{23}}{\partial x_2} + \frac{\partial \sigma_{33}}{\partial x_3} + \gamma_3 = 0$$

$$\tau_{31} = \tau_{13}$$

How many in 2D?

In indicial notation, one can write

$$\sum_{i=1}^3 \frac{\partial \sigma_{ij}(\underline{x})}{\partial x_i} + \gamma_j = 0$$

$$\underline{\underline{\sigma}}(\underline{x}) = \underline{\underline{\sigma}}^T(\underline{x})$$

We have six equilibrium equations, but nine unknown stress components

We need more equations which we will obtain from **stress-strain relation** & **strain-displacement**

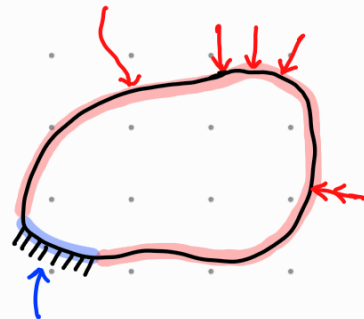
The above equations are partial differential equations, and to solve them **uniquely**, we would need BOUNDARY CONDITIONS

There are two types of boundary conditions

→ displacement BCs (Dirichlet BCs)
specifies displacements on the boundary
(will cover later in deformation)

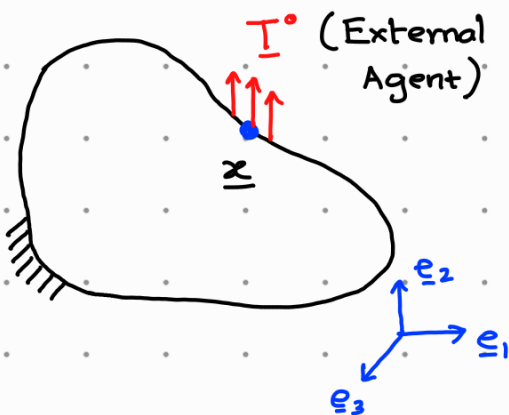
→ traction / force BCs (Neumann BCs)
specifies external loads applied on the boundary

traction BC



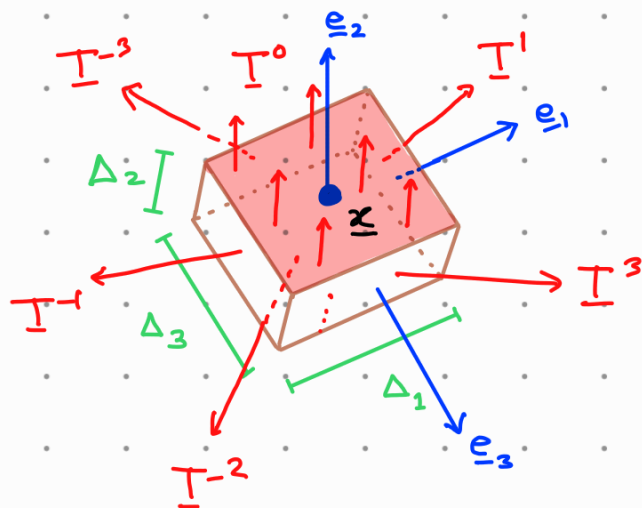
displacement BC (e.g. clamped)

Derivation of traction boundary condition



This applied load is usually distributed over an area; therefore has unit of traction. We denote this distributed $\underline{T}^\circ$

Let's take a small piece from the surface. The actual boundary is curved but if we take a tiny piece, it will be almost flat



On the top surface of this small body part acts the external uniform distributed traction $\underline{T}^0$

The directions of $\underline{e}_1 - \underline{e}_3$ are in the plane of the surface and $\underline{e}_2$ is the normal to the plane.

What are the forces acting on this small body part?

- External force $\underline{T}^0$ (acting on exposed surface)
- Traction (acting on the internal faces)
- Body force (acting at the COM of the volume)

The total force on this small body part must be also be in equilibrium.

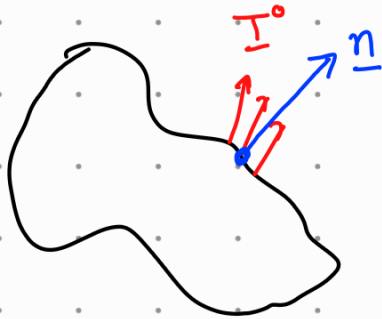
$$\underbrace{\underline{T}^0 \Delta_1 \Delta_3}_{\text{Ext force}} + (\underline{T}^1 + \underline{T}^{-1}) \Delta_2 \Delta_3 + (\underline{T}^3 + \underline{T}^{-3}) \Delta_1 \Delta_2 + \underline{T}^{-2} \Delta_1 \Delta_3 + \underbrace{\gamma \Delta_1 \Delta_2 \Delta_3}_{\text{Body force}} = \underline{0}$$

If you divide the above equation by $\Delta_1 \Delta_3$, we get:

$$\Rightarrow \underline{T}^0 + (\underline{T}^1 + \underline{T}^{-1}) \frac{\cancel{\Delta_2}^0}{\Delta_1} + (\underline{T}^3 + \underline{T}^{-3}) \frac{\cancel{\Delta_2}^0}{\Delta_3} + \underline{T}^{-2} + \gamma \cancel{\Delta_2}^0 = 0$$

$$\boxed{\underline{T}^0 = -\underline{T}^{-2}} \Rightarrow \underline{T}^0 = \underline{T}^2 \Rightarrow \underline{T}^0 = \underline{\underline{\sigma}} \underline{e}_2$$

Now, let $\Delta_2 \rightarrow 0$, that is we shrink the height by pushing the bottom surface towards the top surface while keeping the surface area Δ_1, Δ_3 constant.



$$\underline{\underline{\sigma}} \underline{n} = \underline{I}^o$$

General formula

$$\begin{bmatrix} \sigma_{11} & T_1^o & z_{31} \\ T_1^o & T_2^o & T_3^o \\ z_{31} & T_3^o & \sigma_{33} \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} = \begin{bmatrix} T_1^o \\ T_2^o \\ T_3^o \end{bmatrix}$$

$\underline{\underline{\sigma}}$
 $\underline{e}_2$
 $\underline{I}^o$

