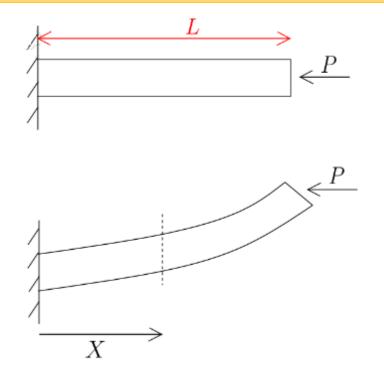
# Classroom Lecture 33

Buckling of beams (contd...) Introduction to Energy methods

# Buckling of beams

#### Beam buckling (contd from last lecture)



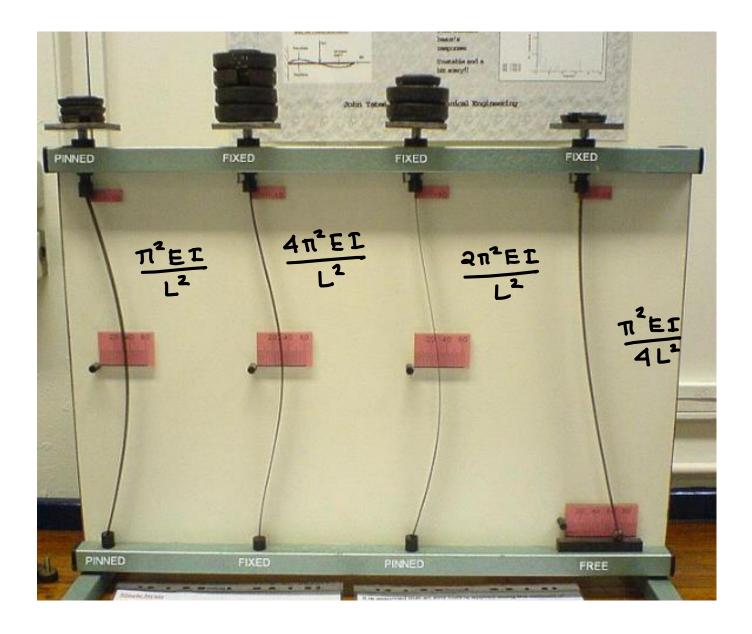
- When the compressive load P reaches the critical value, the beam/column will bend, even though we are applying axial compressive load → Buckling phenomena
- The minimum value of the axial force at which the beam will buckle is the called the critical buckling load,  $P_{cr}$
- For the cantilever beam, *P<sub>cr</sub>* was derived as:

$$P^{cr} = \frac{\pi^2 EI}{4L^2}$$

- If the compressive axial force is less than  $P_{cr}$ , the beam will remain straight
- $P_{cr}$  is proportional to bending stiffness EI and inversely proportional to square of the beam's length
  - Longer beam requires lesser force to buckle

#### Beam buckling depends upon the boundary conditions

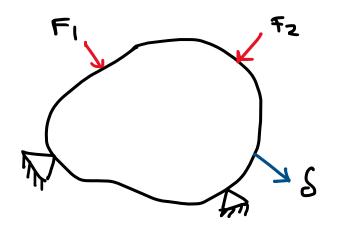
- In class, *P<sub>cr</sub>* was derived for cantilever beam, using the relevant boundary conditions
- The critical buckling load also depends upon the boundary conditions of the beam/column
- Some typical boundary conditions encountered are:
  - Pinned-pinned
  - Fixed-fixed
  - Fixed-pinned
  - Fixed-free (i.e. cantilever)



## Introduction to Energy Methods

#### How did we solve for displacements of a body given applied forces?

• We have solved equations of 3D-elasticity in the form of PDEs to obtain deformation of a body



$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + b_x = \rho a_x$$
$$\frac{\partial \tau_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} + b_y = \rho a_y$$
$$\frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \tau_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} + b_z = \rho a_z$$

 $\underline{\underline{\sigma}} \underline{\underline{n}} = \underline{\underline{t}}_0$ , (traction boundary condition)  $\underline{\underline{u}} = \underline{\underline{u}}_0$ . (displacement boundary condition)

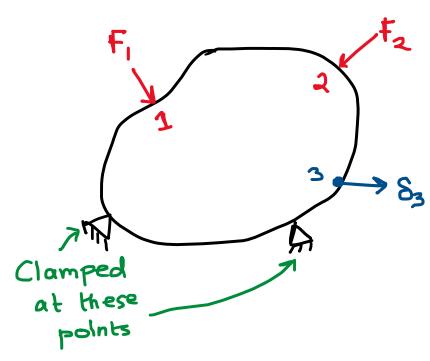
• In last two lectures, we solved ODEs to solve beam deflection problems

$$EI\frac{d^2y}{dX^2} = M(X) \qquad EI\frac{d\theta}{dX} = M(X) \quad \frac{dy}{dX} - \theta = \frac{V}{kGA}$$

• There is also another approach of finding deformation, using energy methods!

Recall that you could solve motion of rigid bodies using either Newton's 2<sup>nd</sup> law or using minimization of energy For many complex problems, it becomes easier to find deformations using energy methods! **Concepts** needed for energy methods

#### How is the displacement at a point related to a force at another point?



- Consider a body subjected to (concentrated) forces at point 1 and point 2, and we are interested in measuring the displacement at point 3
- What is the relationship between the displacement  $\delta_3$  at point 3 and the force  $F_1$  at point 1?

 $\delta_3 \equiv F_1$ 

• Is the relationship linear?

• If linear, what would be proportionality constant?

• To answer these questions, let's have a look at the 3D equations of elasticity

#### Linearity of the elasticity equations and boundary conditions

• 3D equations of elasticity (a set of PDEs)

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + b_x = \rho a_x$$
$$\frac{\partial \tau_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} + b_y = \rho a_y$$
$$\frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \tau_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} + b_z = \rho a_z$$

 $\underline{\underline{\sigma}} \underline{\underline{n}} = \underline{\underline{t}}_0$ , (traction boundary condition)  $\underline{\underline{u}} = \underline{\underline{u}}_0$ . (displacement boundary condition)

• The relation between stress and displacement is linear! How?

 $\blacktriangleright$  These equations are linear in stress components  $\sigma$  and  $\tau$  (e.g. power of stress is one)

- Stress is linear in strain  $\sigma_{ij} = \lambda tr(\underline{\epsilon}) \delta_{ij} + 2\mu \epsilon_{ij}$
- Strain is linear in displacement  $\underline{\boldsymbol{\varrho}} = \frac{1}{2} \left[ \nabla \underline{\boldsymbol{u}} + (\nabla \underline{\boldsymbol{u}})^T \right]$
- Boundary conditions are also linear in displacement  $\underline{u}$

So elasticity equations and the boundary conditions are all linear in the unknown displacement

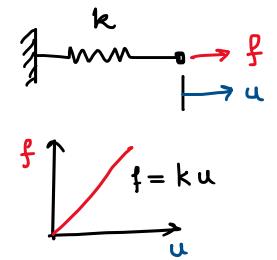
### What does linearity in equations imply?

- Ex 1: Linear spring
  - u is unknown static disp, f is known static force
  - k u = f
  - Double the force *f* , you get double *u*



- If you apply  $f_1(t) \rightarrow$  you get  $u_1(t)$
- If you apply  $f_2(t) \rightarrow you \text{ get } u_2(t)$
- If you apply  $f_1(t) + f_2(t) \rightarrow$  you get  $u_1(t) + u_2(t)$

Superposition principle holds true for linear systems!



#### Linearity and superposition apply to 3D linear elasticity equations as well

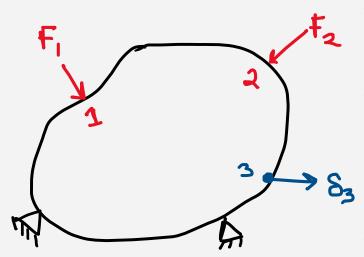
• Here, we have a set of PDEs, and they are still linear!

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + b_x = \rho a_x$$
$$\frac{\partial \tau_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} + b_y = \rho a_y$$
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 $\underline{\underline{\sigma}} \underline{\underline{n}} = \underline{\underline{t}}_0$ , (traction boundary condition)  $\underline{\underline{u}} = \underline{\underline{u}}_0$ . (displacement boundary condition)

- What the forces that act here?
  - Body forces:  $b_x$ ,  $b_y$ ,  $b_z$
  - Surface tractions: <u>to</u> (through boundary conditions)
- If you double these forces, the resulting deformation will also get doubled!

#### How is the displacement at a point related to a force at another point?



- Consider a body subjected to (concentrated) forces at point 1 and point 2, and we are interested in measuring the displacement at point 3
- What is the relationship between the displacement  $\delta_3$  at point 3 and the force  $F_1$  at point 1?

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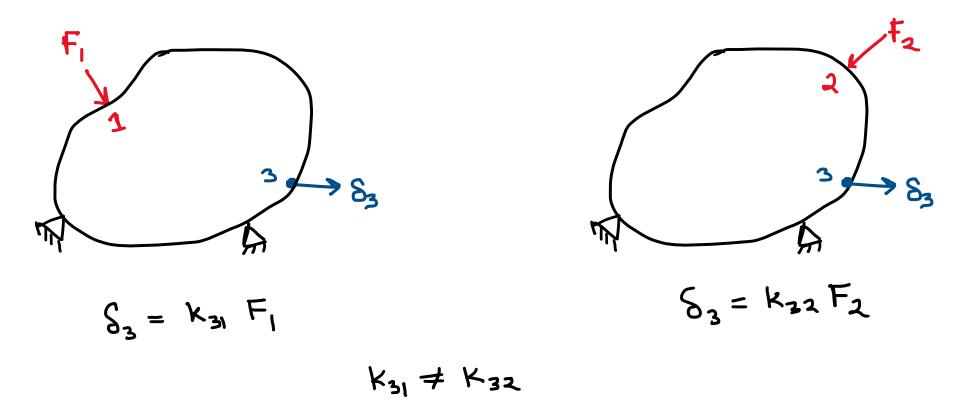
• Is the relationship linear?

- If linear, what would be proportionality constant?
- Since the unknown displacement is **linearly related** to the force, we can write

$$\delta_3 \propto F_1 \; \Rightarrow \; \delta_3 = k_{31}F_1$$

•  $k_{31}$  is the proportionality constant and is called the **influence coefficient** 

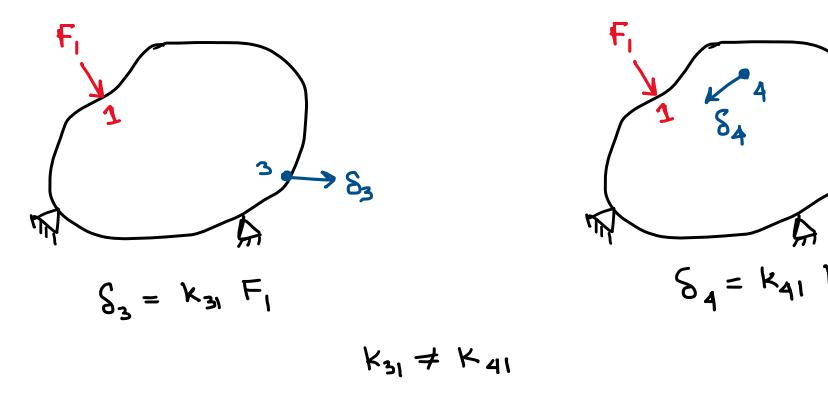
• Influence coefficients vary depending upon the location of the applied force



Influence coefficients depend on

• The location of the applied force

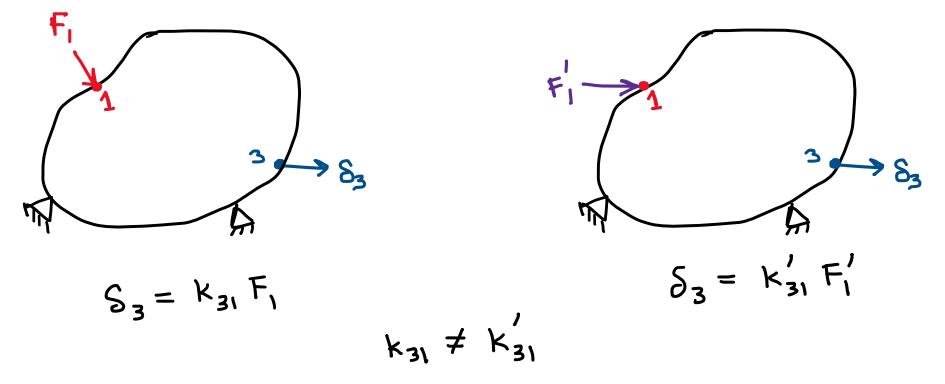
• Influence coefficients vary depending upon the location of the measured displacement



Influence coefficients depend on

• The location of the applied force and the location of the the measured displacement

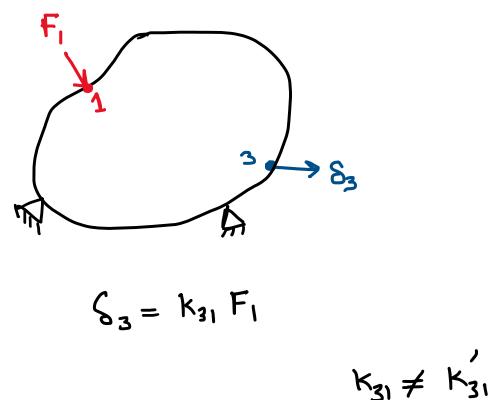
• Influence coefficients depend on the direction of the applied force

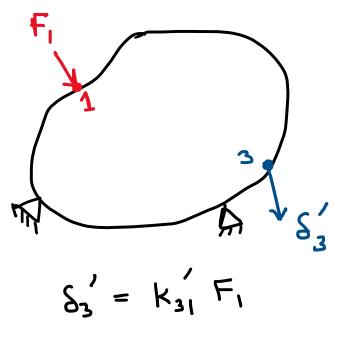


#### Influence coefficients depend on

- The location of the applied force and the location of the the measured displacement
- The direction of the applied force but not on the value of the force!

• Influence coefficients depend on the direction of the measured displacement

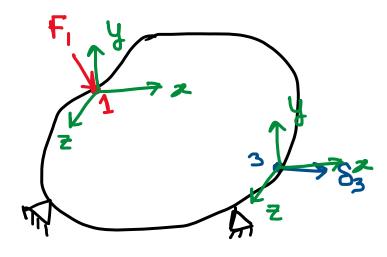




#### Influence coefficients depend on

- The location of the applied force and the location of the the measured displacement
- The direction of the applied force but not on the value of the force!
- The direction in which the displacement is measured

- Influence coefficients are direction-dependent
- Both displacement and force are vector quantities
- Each have three components in three mutually perpendicular directions



irections  

$$F_{1} = \begin{cases} F_{1}^{x} \\ F_{1}^{y} \\ F_{1}^{z} \\ F_{1}^{z} \end{cases}$$
Component  $f$   
in  $x$ -direction  

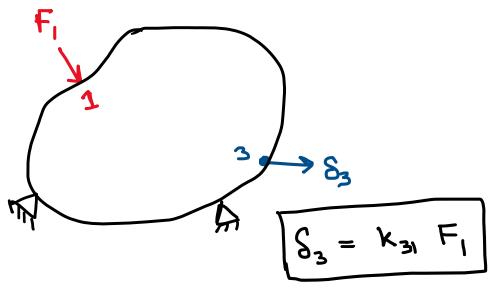
$$F_{1} = \begin{cases} F_{1}^{x} \\ F_{1}^{y} \\ F_{1}^{z} \\ F_{1}^{z}$$

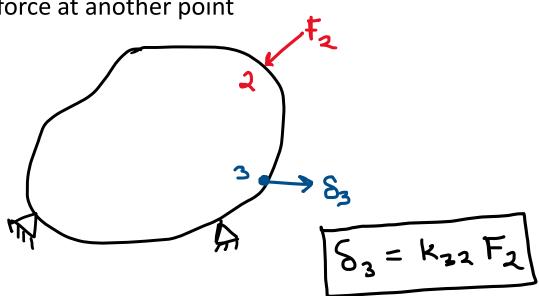
$$\begin{aligned} \mathbf{F}_{3}^{x} &= k_{31}^{xx} F_{1}^{x} & \delta_{3}^{y} = k_{31}^{yx} F_{1}^{x} & \delta_{3}^{z} = k_{31}^{zx} F_{1}^{x} \\ \delta_{3}^{x} &= k_{31}^{xy} F_{1}^{y} & \delta_{3}^{y} = k_{31}^{yy} F_{1}^{y} & \delta_{3}^{z} = k_{31}^{zy} F_{1}^{y} \\ \delta_{3}^{x} &= k_{31}^{xz} F_{1}^{z} & \delta_{3}^{y} = k_{31}^{yz} F_{1}^{z} & \delta_{3}^{z} = k_{31}^{zz} F_{1}^{z} \end{aligned}$$

Each equation resemble the Hooke's law (where stress was linearly related to strain)

#### Displacement expressed as superposition using influence coefficients

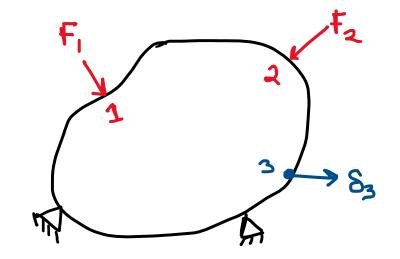
• Displacement at one point is a linear function of applied force at another point



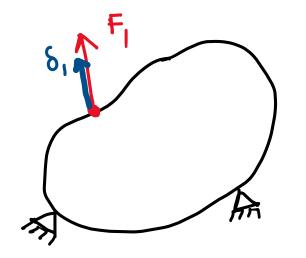


 Due to superposition, displacement at a point due to forces applied at two different points can be written using superposition principle

$$S_3 = k_{31}F_1 + k_{32}F_2$$



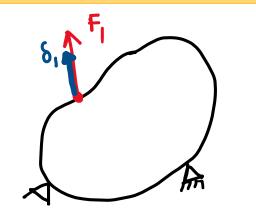
#### Concept of corresponding displacement



- The component of the displacement of a body at the same location and along the same direction as the applied force is called the "corresponding displacement"
- Note that the resulting displacement vector may not point in the direction of the applied force
- The "corresponding displacement" is responsible for actual work done, since the component of displacement perpendicular to force is zero. Hence, it is also called the *work-absorbing displacement*.
- If we apply force  $F_1$  at point 1 and measure the corresponding displacement  $\delta_1$ , then they are related linearly as:

$$\delta_1 = k_{11} F_1$$

### Energy stored in a deformable body due to corresponding displacement



- How much is the energy stored in a deformable body due to an applied force  $F_1$ ?
- Would it be: Stored energy =  $F_1$ .  $\delta_1$ ?